

Uniformly Best Wavenumber Approximations by Spatial Central Difference Operators

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Classical Finite Differences

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Philosophy...

Consider a central finite difference stencil

$$\left. \frac{du}{dx} \right|_{x=x_j} \approx \sum_{k=1}^P c_k^{(p)} (u_{j+k} - u_{j-k})$$

Philosophy

Error in numerical solution is governed by the truncation error of the finite difference stencil.

Strategy

Choose $c_k^{(p)}$ to eliminate terms up to $\mathcal{O}(\Delta x^{2p})$

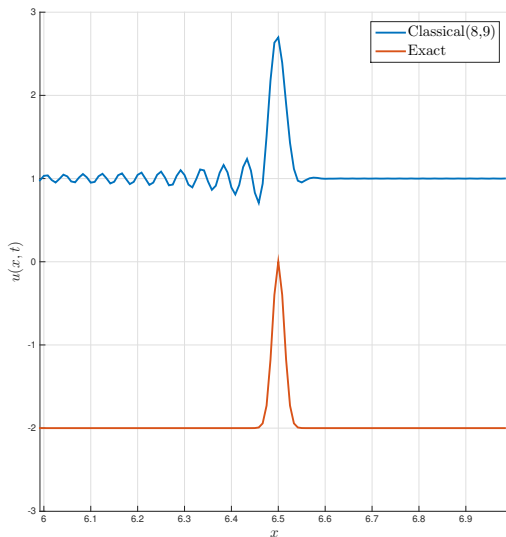


Example - Advection equation

$$u_t + u_x = 0, \quad 0 \leq x < 6, \quad t \geq 0$$

$$u(x, 0) = 2 \exp\left(-3200 \left(x - \frac{1}{2}\right)^2\right)$$

Periodic Boundary Conditions



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Analytic dispersion relation

$$u_t + u_x = 0$$

Assume u has uniformly convergent Fourier series. Consider general Fourier mode, $u(x, t) = \exp(i(\kappa x - \omega t))$

$$\omega = \kappa, \quad \text{Analytic dispersion relation}$$

$$v_p = \frac{\omega}{\kappa}, \quad \text{Phase speed}$$

$$v_g = \frac{d\omega}{d\kappa}, \quad \text{Group speed}$$

Speeds independent of wavenumber

⇒ **Non-dispersive solution!**

Numeric dispersion relation

$$(u_t)_j + \sum_{k=1}^p c_k^{(p)} (u_{j+k} - u_{j-k}) = 0$$

$$u_j(t) = \exp(i(\kappa x_j - \bar{\omega} t))$$

$$\Rightarrow \underbrace{\bar{\omega} \Delta x}_{\bar{\xi}} = 2 \sum_{k=1}^p c_k^{(p)} \sin(k \underbrace{\kappa \Delta x}_{\xi})$$

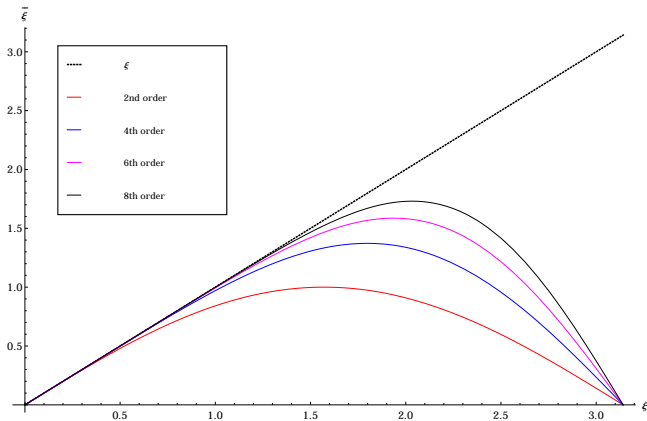
Speeds dependent on wavenumber

\Rightarrow **Inherently dispersive solution!**

Incorrect phase speed, $\bar{v}_p = \frac{\bar{\omega}}{\kappa}$, and group speed, $\bar{v}_g = \frac{d\bar{\omega}}{d\kappa}$

Can show that $\bar{\xi} < \xi$, $\bar{v}_p < v_p$, $\bar{v}_g < v_g$ for all classical stencils.
In fact

$$\xi - \bar{\xi} = \frac{2^p}{\binom{2p}{p}} \int_0^\xi (1 - \cos(\xi'))^p d\xi'$$



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Alternative philosophy

Problems involving high frequencies / wavenumbers that travel for long times have errors dominated by bad numerical dispersion.

- Fluid dynamics
- Aeroacoustics
- Electromagnetism
- Seismology

Philosophy

Consider dispersion error

$$E(\xi, \mathbf{a}) = \xi - \bar{\xi}(\xi, \mathbf{a})$$

when choosing stencil coefficients

A new problem...

Consider finite difference stencil

$$\left. \frac{du}{dx} \right|_{x=x_j} = \sum_{k=1}^{p+n} a_k (u_{j+k} - u_{j-k}) + \mathcal{O}(\Delta x^{2p})$$

$$\bar{\xi} = 2 \sum_{k=1}^{p+n} a_k \sin(k\xi)$$

- Accuracy constraint - Order $\mathcal{O}(\Delta x^{2p})$
- Leaves n degrees of freedom
- Use these to minimise dispersion error uniformly, i.e. $\|\xi - \bar{\xi}\|_\infty$

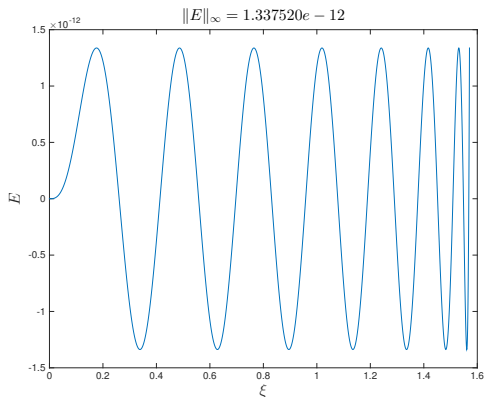


Main results

Theorem

There is a unique stencil \mathbf{a} that minimises $\|\xi - \bar{\xi}(\xi, \mathbf{a})\|_\infty$. The error of this stencil oscillates $n + 1$ times.

Can devise a convergent algorithm for finding best possible \mathbf{a}
(Remez algorithm)



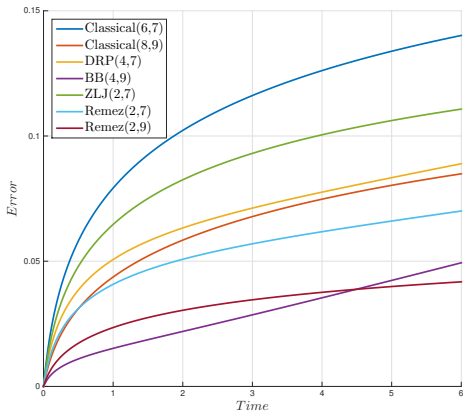
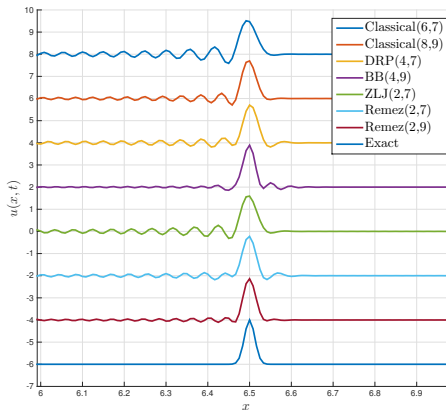


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Periodic Boundary Conditions



- [1] Tam, Webb (1993)
- [2] Zingg, Lomax, Jurgens (1996)
- [3] Bogey, Bailly (2004)

Conclusion

- For many problems the numerical error comes from inaccurate approximation of Dispersion Relation
- Classical philosophy does not account for this
- We can construct accurate Finite Difference stencils with arbitrarily small Dispersion Error

For more information, see

Linders, Nordström, Journal of Computational Physics, 2015